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In situ spectroscopic ellipsometry during electrochemical treatment of zinc in alkaline carbonate electrolyte

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Abstract

In situ spectroscopic ellipsometry (SE, 297–820 nm, 4.4–1.5 eV) during electrochemical oxidation/reduction of zinc in alkaline carbonate electrolyte was performed using a specially designed optical-electrochemical cell. The growth and shrinkage of the interfacial layer on Zn was analysed on the basis of Lekner's first order perturbation theory. For non-absorbing thin films on metal surfaces, an algorithm has been derived to extract the thickness of a surface film directly from ellipsometric data without the need of optical modelling. During cyclic voltammetry (CV), a rapid decrease/increase in the layer thickness in the reduction/oxidation peak has been found. In the potential regime where the surface is oxide-covered, the layer thickness increases/decreases linearly with potential in anodic/cathodic scans. The density of the interfacial region in this regime is constant. During chronoam-

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¹Abbreviations used: CA – chronoamperometry, CV – cyclic voltammetry, OCP – Open Circuit Potential, PCSA – Polariser-Compensator-Sample-Analyser, SE – spectroscopic ellipsometry, SHE – Standard Hydrogen Electrode

perometric experiments, a fast correlating jump in thickness was found after potential jumps. An ageing of the films is observed in the absorption spectrum (from the ellipsometric parameter Ψ), which shows changes until ≈ 30 min after potential jumps. Analysis of the current transients points to continuous dissolution of Zn.

Keywords: ellipsometric perturbation analysis, aqueous corrosion, passivity, zinc hydroxide, zinc oxide, semiconductor thin film

1. Introduction

Zn is one of the most important metals in applications. Due to its oxidation propensity, Zn is a very efficient sacrificial anode in cathodic protection, and used as metallic coating [1]. For most applications, as well as in the deposition process of metallic coatings, electrochemical reactions on the Zn surface are of great importance. In recent years, ZnO as a semiconducting material with wide band gap of 3.4 eV has generated interest due its electronic properties [1–4]. ZnO is commonly found in the corrosion products of Zn and its alloys [1, 3]. The semiconducting properties of ZnO have been found to play an important role in the corrosion behaviour of Zn [1, 5–7]. Electrochemical works have shown that oxide films with different properties form in different electrolytes, some passivating the zinc surface better than others [1, 8].

A number of works focused on the electrochemistry of zinc over a range of pH values. The passivation of zinc, typically in alkaline media, has been investigated, and different kinds of passivating films have been proposed [1, 8–12]. Many passive films are amorphous or nanocrystalline [13, 14], with

18 broad distribution functions of distances and coordination numbers instead
19 of sharp, well-defined values of crystals [3, 15, 16].

20 The complexity of the chemical – in addition to electrochemical – reac-
21 tions on the Zn surface, in particular at alkaline pH, implies the need for
22 additional surface characterisation, preferably in situ, as many of the po-
23 tentially formed species are prone to structural changes during the drying
24 process needed for modern surface analytics in UHV [1, 3, 16].

25 Ellipsometry in general is a sensitive optical method for the measurement
26 of the optical constants and thicknesses of thin films [10, 17–19]. In situ
27 ellipsometry can be performed during the modification process of a sample,
28 which can be, e.g. the growth of a thin film, etching or cleaning of a sample
29 [10, 20–23]. The technique enables the determination of fundamental process
30 parameters, such as growth or etch rates and variation with time of optical
31 properties, the latter related to structural changes. Surface reactions on
32 many metals and metallic materials, such as iron, copper, steel, aluminium,
33 zinc and silver have been studied by this technique [2, 9, 21, 24, 25]. The
34 thickness results obtained by ellipsometry, frequently as a result of fitting
35 to an optical model, have been used to evaluate the growth mechanism and
36 kinetics of oxide films on different metals [2, 9, 24, 25]. One result important
37 for Zn is that after reduction for more than 10 min at a potential $E = -1.4$
38 V vs. saturated calomel electrode, a Zn surface without the presence of any
39 oxide can be produced [9].

40 Ellipsometric measurements can either be performed at a single wave-
41 length, or in a spectroscopic ellipsometry (SE) experiment at multiple wave-
42 lengths. In the latter case, the recording of spectral changes provides in-

43 formation on the interfacial electronic structure. SE can also be coupled to
44 electrochemical experiments to study electrochemical reactions. The electro-
45 chemical activation and oxide growth on glassy carbon electrode was in situ
46 measured by SE combined with other techniques [26]. The electrochemical
47 oxidation of nickel and precipitation of hydroxide films was observed by in
48 situ SE [27]. Recently, in situ SE has been used to monitor the transient
49 growth and thinning of the passive film on Fe in borate buffer [20]. More-
50 over, the growth of oxide on a Zn substrate in different atmospheres has been
51 investigated with in situ SE [28].

52 In this work, SE was used for in situ investigations of electrochemical pro-
53 cesses on the Zn/electrolyte interface. Two different electrochemical experi-
54 ments were performed to study the transformations on the surface, (a) cyclic
55 voltammetry (CV) with dynamic potential scans and (b) chronoamperome-
56 try (CA) at certain static potentials, both in carbonate solution. Carbonate
57 electrolyte was used as the complete exclusion of carbonate is difficult under
58 ambient conditions, but affects the reactivity. An analysis based on the per-
59 turbation parameter of Lekner [29] is introduced and used to analyse the SE
60 data.

61 **2. Materials and Methods**

62 *2.1. Materials and sample preparation*

63 Zn foil (99.95%, Goodfellow, Germany) with a thickness of 2 mm was cut
64 into pieces of 20 mm \times 30 mm. The sample was then mechanically ground
65 with SiC paper up to 2500 grit, and subsequently further polished with final
66 polishing suspension (100 nm SiO₂, pH = 9) to minimise surface roughness.

67 After that, these samples were sonicated in ethanol for 10 min and dried in
68 a nitrogen stream. To reduce possible contamination and further oxidation,
69 the sample after this treatment was directly put in the cell for electrochemical
70 measurement. As the surface is already oxide-covered and was exposed to
71 carbonate, surface reactions on the time-scale of the transfer are expected to
72 have only a limited effect on the layer. The oxide is protecting the Zn, and
73 after transfer is immersed in a solution similar to the previous one, so that
74 the oxide/electrolyte interface is expected to be in a similar state as before
75 the transfer. Therefore, no special transfer chamber to avoid contact with
76 the outside air was constructed.

77 *2.2. Design of optical-electrochemical cell*

78 Details of the design of the sample cell were described in a different con-
79 text elsewhere [28]. In addition to the previously described application of
80 measurements in different gas atmospheres, the cell can also be filled and
81 rinsed with different solutions. Four optical windows enable measurements
82 at two incident angles, 50° and 70°. The cell contains a conventional three-
83 electrode electrochemical setup. A platinum mesh as counter electrode and
84 an Ag/AgCl/3M KCl electrode as reference electrode (DriRef-2SH microref-
85 erence electrode, $E = +0.208$ mV vs. standard hydrogen electrode (SHE),
86 World Precision Instruments Inc., USA) were used for electrochemical treat-
87 ment. The connection of a metallic sample was realised by putting it in direct
88 contact with the copper base part, which was connected to the working elec-
89 trode connection of the potentiostat. In this work, all measurements were
90 done at an incidence angle of 70°.

91 *2.3. Electrochemistry*

92 Electrochemical treatment of Zn was performed in 1.0 M Na₂CO₃ aqueous
93 solution with a Gamry potentiostat (PHE200, Gamry Instruments Inc., USA)
94 at room temperature ($22 \pm 2^\circ\text{C}$). Most data presented here were obtained in
95 the cell described above. In this cell, the electrode surface area is well-
96 defined, which is why the results of these measurements are presented as
97 current densities I .

98 For reference, additional CV measurements were conducted in a standard
99 electrochemical cell in a conventional three-electrode setup, where the current
100 is flowing between the working electrode (i.e. the surface under study) and a
101 counter electrode, while the reference electrode is used to precisely measure
102 the electrode potential difference to the working electrode [30]. In this cell,
103 the surface area A was difficult to determine because the electrode was Zn
104 sheet with two sides and also the faces and corners exposed to the electrolyte.
105 Therefore, for this cell (only in Fig. 1) current i instead of I is given.

106 *2.4. Ellipsometry*

107 In-situ SE measurements were performed using a UV-Vis spectroscopic
108 ellipsometer (SE 800, Sentech Instruments GmbH, Germany) working in the
109 wavelength range 297 nm – 820 nm (4.4 – 1.5 eV). In each measurement, data
110 for 856 wavelengths were recorded on the instrument’s CCD line detector.
111 The light source was a Xenon lamp. For in situ measurements, the sample
112 was mounted inside the optical-electrochemical cell described above.

113 In PCSA geometry [17], the polariser was fixed at an angle of $+45^\circ$.
114 Intensity measurements at 20 different analyser positions were used with and
115 without a fixed retarder using the protocols implement in the instrument’s

116 control software. For each wavelength λ in the range between 297 to 820 nm,
 117 the ellipsometric angles Ψ and Δ were extracted [31]. Ellipsometric spectra
 118 were acquired during CVs with a scan rate of 2 mV s⁻¹ between +0.3 V
 119 and -1.7 V where acquisition of one spectrum took (43 ± 1) s, as well as
 120 during CA (-1.3, +0.2 and +1.2 V) measurements, where acquisition of one
 121 spectrum took (34 ± 1) s.

122 The angles Ψ and Δ were transformed into the ellipsometric ratio $\rho =$
 123 $\frac{r_p}{r_s} = \tan(\Psi) \exp(j\Delta)$ with $j = \sqrt{-1}$ [17]. This ratio of the amplitude re-
 124 flection coefficients r_p and r_s for p- and s-polarised light, respectively, was
 125 analysed using a perturbation approach [29]. For individual layers and layer
 126 systems with a total thickness small compared to the wavelength, ρ is ex-
 127 panded to first order around a step profile in the dielectric function with ρ_0 ,
 128 yielding [29]

$$\rho = \rho_0 - \frac{2jq_1K^2(q_1 + q_2)}{(q_1 - q_2)\epsilon_1^2\epsilon_2(q_1/\epsilon_1 + q_2/\epsilon_2)^2}J_1, \quad (1)$$

129 where the amplitude reflection coefficient $r_{s,0}$ of the step profile was replaced
 130 by the respective Fresnel equation [29]. Here,

$$K = \frac{2\pi\sqrt{\epsilon_1}}{\lambda} \sin(\theta_1) \quad (2)$$

131 is the wavevector component parallel to the interface at which a plane wave
 132 impinges under an angle of θ_1 , while

$$q_o = \frac{2\pi\sqrt{\epsilon_o}}{\lambda} \cos(\theta_o), o \in 1, 2 \quad (3)$$

133 is the wavevector component perpendicular to the interface in the respec-
 134 tive medium, where index 1 indicates the medium of incidence and 2 indi-
 135 cates the Zn substrate. Literature values have been used for the wavelength-
 136 dependence of the dielectric function ϵ_1 of the electrolyte [32]. For the Zn

137 substrate, the Drude-Lorentz model with parameters previously reported has
 138 been used [7]. Eq. 1 can be solved for J_1 and used for a computation of J_1
 139 from experimental data.

140 The perturbation parameter J_1 is related to the transition of the dielectric
 141 function of the interfacial region $\epsilon_l(z, \lambda)$ in z -direction (which is perpendicular
 142 to the interface) [29],

$$J_1(\lambda) = \int_{-\infty}^{\infty} \frac{[\epsilon_1(\lambda) - \epsilon_l(z, \lambda)][\epsilon_l(z, \lambda) - \epsilon_2(\lambda)]}{\epsilon_l(z, \lambda)} dz. \quad (4)$$

143 A measurement of the bare Zn/electrolyte interface without interfacial
 144 film may be difficult because of the reactivity of Zn [1, 33]. However, at
 145 cathodic potentials, complete removal of the surface oxide film has been
 146 reported [9], a result which is confirmed here. Therefore, ρ at cathodic po-
 147 tentials has been used as ρ_0 . During CV, data for the potential range -1.4
 148 V to -1.2 V has been averaged to obtain ρ_0 . During CA, two measurements
 149 at -1.3 V have been averaged to calculate ρ_0 .

150 2.5. Analysis of J_1

151 In the following, consider a homogeneous interfacial layer on top of the
 152 Zn substrate with an effective thickness d and an effective dielectric function
 153 $\epsilon_l(\lambda)$. For this case, eq. 4 simplifies to [29]

$$J_1(\lambda) = \frac{[\epsilon_1(\lambda) - \epsilon_l(\lambda)][\epsilon_l(\lambda) - \epsilon_2(\lambda)]}{\epsilon_l(\lambda)} d. \quad (5)$$

154 Previously, the case of an absorbing layer on a non-absorbing substrate
 155 was treated in detail [29]. Further, for non-absorbing, anisotropic systems,
 156 first order [34, 35] and second order perturbation parameter [35] have been
 157 used to analyse interfacial properties of liquid crystals. Zinc oxide, hydroxide

158 and carbonate species to be investigated here do typically not absorb light
 159 at wavelengths in the upper part of the visible spectrum, therefore ϵ_l can
 160 be treated as purely real, i.e. $\epsilon_l = \text{RE}(\epsilon_l)$ [7, 36]. Working on a metallic
 161 substrate, J_1 will nevertheless have an imaginary part. Therefore, here the
 162 case of a non-absorbing film on an absorbing substrate shall be considered.

163 The electrolyte is also non-absorbing, i.e. $\epsilon_1 = \text{RE}(\epsilon_1)$. Further, the
 164 layer thickness is real, $d = \text{RE}(d)$. Therefore, from eq. 5, two equations are
 165 obtained of real variables,

$$\text{RE}(J_1) = \left[\epsilon_1 + \text{RE}(\epsilon_2) - \epsilon_l - \frac{\epsilon_1}{\epsilon_l} \text{RE}(\epsilon_2) \right] d \quad (6)$$

166 and

$$\text{IM}(J_1) = \left(1 - \frac{\epsilon_1}{\epsilon_l} \right) \text{IM}(\epsilon_2) d, \quad (7)$$

167 which can be solved for the two unknown quantities d and ϵ_l . (From the
 168 purely mathematical point of view, a similar approach can be used to invert
 169 ρ directly to d and ϵ_l in a ambient/layer/substrate system. However, a much
 170 higher precision in measurement is needed to yield stable results using such
 171 an approach compared to the perturbation approach.)

172 Taking the ratio of eq. 6 and 7 removes the dependence on d ,

$$\frac{\text{RE}(J_1)}{\text{IM}(J_1)} = \frac{\epsilon_1 + \text{RE}(\epsilon_2) - \epsilon_l - \frac{\epsilon_1}{\epsilon_l} \text{RE}(\epsilon_2)}{\left(1 - \frac{\epsilon_1}{\epsilon_l} \right) \text{IM}(\epsilon_2)}. \quad (8)$$

173 Solving for ϵ_l yields

$$\epsilon_l = \text{RE}(\epsilon_2) - \text{IM}(\epsilon_2) \cdot \frac{\text{RE}(J_1)}{\text{IM}(J_1)}. \quad (9)$$

174 This result is subsequently substituted into eq. 6 or 7, which have then been
 175 solved for d . The structure of eq. 9 shows that this procedure is not applicable

176 for sample systems containing only non-absorbing materials, as in this case
177 $\text{Im}(J_1) = 0$.

178 While ϵ_l depends on the wavelength, d does not. One ellipsometric spec-
179 trum contains ρ at many wavelengths. Here, typically results for d obtained
180 numerically from the procedure described for $\lambda > 650$ nm were averaged to
181 yield d as reported in section 3. While for each value of d , 272 values at
182 different wavelengths have been averaged, ϵ_l remains a function of the wave-
183 length, so one point in the spectrum yields only a single ϵ_l . The data quality
184 is not sufficient to analyse the spectral changes directly from J_1 . Rather,
185 information about the changes in the absorption spectrum are directly ob-
186 tained from Ψ . Applying the analysis procedure described here to simulated
187 data for a system Zn/ZnO/Air with different layer thicknesses of ZnO shows
188 that the layer thickness obtained is typically 10-20 % too low, which is still
189 considered as a physically reasonable result. Consequently, the obtained ϵ_l
190 are too high, and sometimes take unphysically high values, which is a further
191 reason for not discussing experimentally obtained values of ϵ_l .

192 The analysis here remains valid if a hydroxide layer instead of an oxide
193 layer is present on the surface, and also if an hydroxide-terminated oxide
194 layer is present. What limits the overall results is how close one can get
195 experimentally to the “ideal” step profile in the refractive index. If, in the
196 reference state, a hydroxidic termination is present on the surface, d can be
197 interpreted as the increase in layer thickness with respect to the reference
198 state. Effects of the ion distribution around the interface, though they can
199 in principle be detected in ellipsometric experiments, are negligible due to
200 the low refractive index contrast compared to the oxidic films.

201 **3. Results and Discussion**

202 *3.1. Electrochemical behaviour of Zn in carbonate*

203 Fig. 1 shows a typical CV of Zn in 1.0 M Na₂CO₃. Two anodic current
204 peaks (A₁ and A₂), one anodic current plateau and two cathodic peaks (C₁
205 and C₂) were observed, consistent with [11]. According to [11], the oxida-
206 tion current ($\text{Zn} \rightarrow \text{Zn}^{2+} + 2 \text{e}^-$) on the bare metal is responsible for the
207 oxidation peak A₁. A residual current passing through the already formed
208 oxide/hydroxide film results in the current peak A₂, indicating that the ox-
209 ide films formed in alkaline carbonate solution are rather porous. As Fig. 1
210 shows, the interfacial film is reduced at about -1.3 V. The reduction peaks
211 C₁ and C₂ originate from a reduction current for ZnO and/or Zn(OH)₂ and
212 the redeposition current for Zn²⁺, respectively. At lower potentials, active
213 surface was present and H₂ evolution is observed.

214 *3.2. In situ SE during CV*

215 Ellipsometric measurements have been performed to determine the change
216 of the interfacial structure during reduction and subsequent oxidation. Dur-
217 ing a dynamic potential scan as in a CV measurement, the potential range
218 over which a single SE spectrum averages depends on the scan rate. If the
219 same scan rate as in Fig. 1 were applied during in situ SE measurements, the
220 potential range for each SE spectrum will be too wide ($50 \text{ mV/s} \cdot 43 \text{ s} = 2.15$
221 V) and the results can not be used to probe potential-dependent surface fea-
222 tures. Thus, a scan rate of 2 mV/s was used. At 2 mV/s, each SE spectrum
223 averages over ≈ 85 mV.

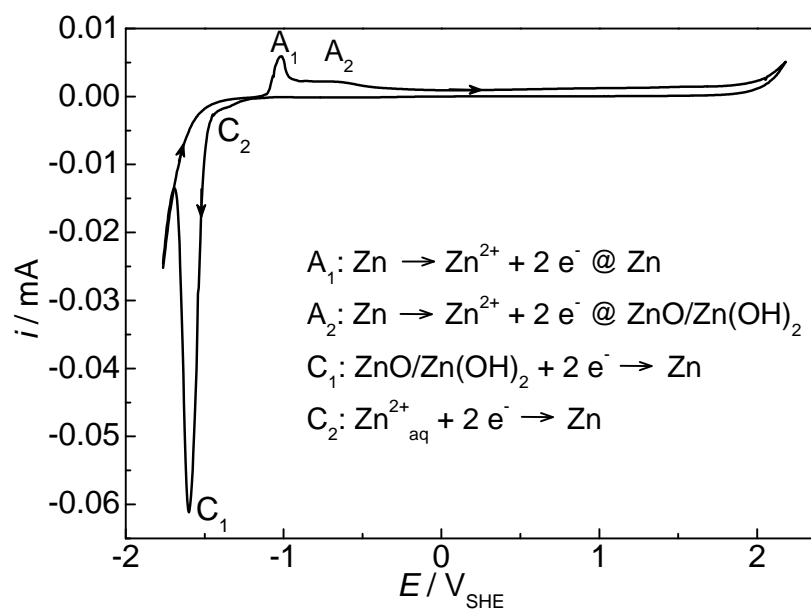


Figure 1: CV of polycrystalline Zn in 1.0 M Na_2CO_3 in electrochemical cell. Scan rate: 50 mV/s. Peak assignments according to [11].

224 Fig. 2 shows the CV at a scan rate of 2 mV/s and the thickness d of
225 the interface film (calculated following the procedure outlined in section 2.5)
226 as function of potential during a CV scan. Before cycling the potential,
227 one SE measurement was performed at open circuit potential (OCP), i.e.
228 at the potential which spontaneously forms when putting into contact an
229 electrode and an electrolyte (empty square in Fig. 2). The OCP itself has
230 not been measured, and the point has been arbitrarily placed on the potential
231 axis. Subsequently, E was changed to the the positive limit, +0.3 V, and
232 a CV measurement was started. At this potential, Zn is covered with an
233 oxidic layer. The film thickness at open circuit potential (OCP) is about
234 0.8 nm. With the application of an anodic potential, the thickness jumps
235 to a value of ≈ 2.2 nm on a time scale below the experimental resolution.
236 During scanning in cathodic direction, a gradual decrease of the thickness
237 was observed till -0.96 V. Upon further reduction, the thickness decreases
238 at a rate of ≈ 8 nm/V, it reaches almost zero at a potential of -1.13 V
239 and keeps constant till -1.48 V. In this potential range, as displayed in
240 the CV in Fig. 2, two cathodic peaks are observed. The cathodic peaks
241 indicate the reduction of oxidised species, resulting in a substantial thinning
242 of the oxide film. All of these features can be directly observed by in situ
243 SE. In the CV, one more broad reduction peak (-1.20 V to -1.45 V) was
244 observed after the first reduction peak (-1.13 V) which does coincide with
245 a region of constant thickness of the interfacial layer. This behaviour was
246 attributed to the fact that the first cathodic peak is due to the reduction of
247 previously formed oxide/hydroxide, while the second reduction peak was due
248 to the redeposition of Zn from previously dissolved Zn^{2+} ions present in the

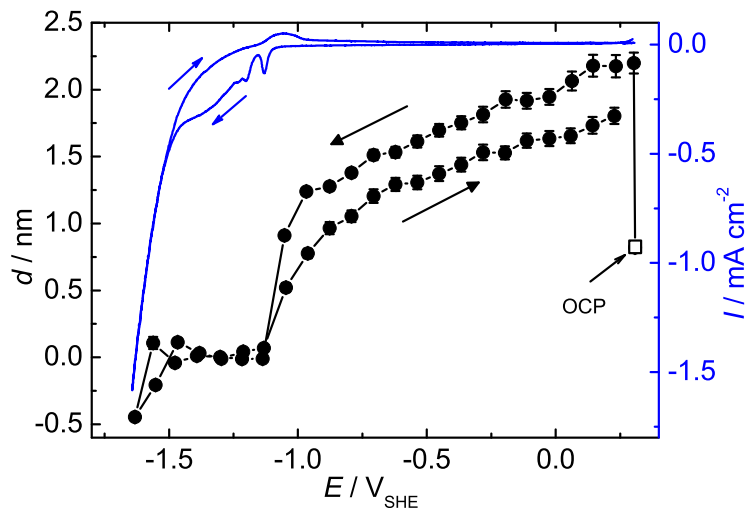


Figure 2: Calculated thickness (■, scale on the left) during a CV (line, scale on the right) at a scan rate of 2 mV/s. The open square (□) is an initial measurement at OCP without electrodes connected.

249 solution in this case. Thus, a cathodic current peak was observed without
 250 the further change of oxide layer thickness between -1.20 V to -1.45 V. This
 251 result confirms previous conclusions from purely electrochemical work about
 252 the current peaks in the CV [11, 12]. From -1.48 to -1.70 V, there are jumps
 253 in the value of the thickness, attributed to the starting H_2 evolution reaction
 254 and the consequent presence of adsorbed H_2 , which leads to formation of
 255 macroscopic bubbles. At the same time, a fast increase of the cathodic
 256 current is observed in CV (Fig. 2).

257 After reversal of the scan direction, no obvious change is seen for the
 258 thickness till the potential reaches -1.13 V, where the oxidation peak A_1
 259 starts. Upon further positive scanning, the thickness increases with ≈ 5 nm/V

260 between -1.13 V and -0.96 V (where peak A_1 ends). At higher potentials,
 261 a drop in current to a plateau value is observed. During this plateau, the
 262 thickness increases almost linearly with potential at 1.0 nm/V. Such a linear
 263 growth with increasing potential was found for surface layers on metals in the
 264 passive region of the current-potential diagram [20], though other behaviour
 265 was reported for Zn in borate buffer [2]. Linear increase of layer thickness in
 266 the passive regime is consistent with what is called the “point defect model”
 267 of passive films [37, 38]. From the electrochemical point of view, the oxide
 268 film formed here is not too “passive”, as the residual current density is $3 - 10$
 269 $\mu\text{A cm}^{-2}$.

270 The anodic current measured during the CV can be used to evaluate the
 271 net integral amount n_{Zn} of Zn over time t which is oxidised to Zn^{2+} , assuming
 272 that only Zn oxidation following the reaction $\text{Zn} \rightarrow \text{Zn}^{2+} + 2 e^-$ contributes
 273 to the current i as

$$n_{\text{Zn}}(E) = \frac{1}{2eN_A} \int_{E(i=0)}^E i(t)dt, \quad (10)$$

274 with Avogadro’s constant N_A and the elementary charge e . The resulting
 275 dependence of n_{Zn} on potential is shown in Fig. 3a. The oxidised Zn^{2+} is
 276 expected to be distributed mainly to three different species: dissolved $\text{Zn}_{\text{aq}}^{2+}$,
 277 $\text{Zn}(\text{OH})_2$ and ZnO . Detailed ex situ analysis of the resulting layers, which is
 278 not topic of this work, shows that only traces of carbonate are present on the
 279 surface, which is why carbonate species are not considered for the analysis.

280 From n_{Zn} and d , an effective concentration c_{Zn} of Zn^{2+} in the layer is
 281 obtained (A - electrode surface area),

$$c_{\text{Zn}} = \frac{n_{\text{Zn}}}{Ad}. \quad (11)$$

282 The potential/time dependence of d was obtained by linear interpolation

283 between the points at which measurements were performed. Fig. 3b shows
 284 c_{Zn} divided by $c_{\text{ZnO}} = 1/V_M(\text{ZnO})$, the reciprocal molar volume V_M of crys-
 285 talline ZnO, $c_{\text{ZnO}} = 6.9 \cdot 10^4 \text{ mol/m}^3$ [32]. The ratio $c_{\text{Zn(OH)}_2}/c_{\text{ZnO}}$ with
 286 $c_{\text{Zn(OH)}_2} = 3.1 \cdot 10^4 \text{ mol/m}^3$ [32] is shown for reference as a line. Differ-
 287 ent values of $c_{\text{Zn}}/c_{\text{ZnO}}$ can be understood considering $c_{\text{Zn}}/c_{\text{ZnO}} \propto Q/d$ with
 288 the net transferred charge Q . If only metal dissolution without formation
 289 of a layer is going on, $d \rightarrow 0$, hence $c_{\text{Zn}}/c_{\text{ZnO}} \rightarrow \infty$. If a layer of crys-
 290 talline ZnO is formed, $c_{\text{Zn}}/c_{\text{ZnO}} = 1$. For amorphous ZnO films, values
 291 slightly lower than 1 are expected. If the product of oxidation is Zn(OH)_2 ,
 292 $c_{\text{Zn}}/c_{\text{ZnO}} = c_{\text{Zn(OH)}_2}/c_{\text{ZnO}} \approx 0.45$. Cases with strongly swollen networks result
 293 in $c_{\text{Zn}}/c_{\text{ZnO}} < c_{\text{Zn(OH)}_2}/c_{\text{ZnO}}$. In addition, $c_{\text{Zn}}/c_{\text{ZnO}} \ll 1$ is also expected if
 294 the layer grows by precipitation from solution rather than electrochemically.

295 The graph in Fig. 3b rises sharply from zero to ≈ 3.3 . The sharp increase,
 296 which goes already beyond the value of 1.0 required for direct ZnO formation,
 297 shows that initially, the current increase results in metal dissolution. In
 298 the region of linear thickness growth, $c_{\text{Zn}}/c_{\text{ZnO}}$ rises only marginally from
 299 3.3 to 4.0. The current flowing in this potential regime therefore mainly
 300 results in a thickening of the interfacial film, and only to a lower extend in
 301 metal dissolution. More information of the nature of the surface film can be
 302 obtained by analysing the spectra.

303 Fig. 4 shows the real and imaginary part of J_1 as a function of potential
 304 for different wavelengths during the CV scan. For all wavelengths, $\text{Re}(J_1)$ be-
 305 haves the same as that of the calculated oxide thickness which was extracted
 306 from high wavelength data only. This fact shows that the assumption of small
 307 changes in the refractive index in the oxide layer is justified. The behaviour

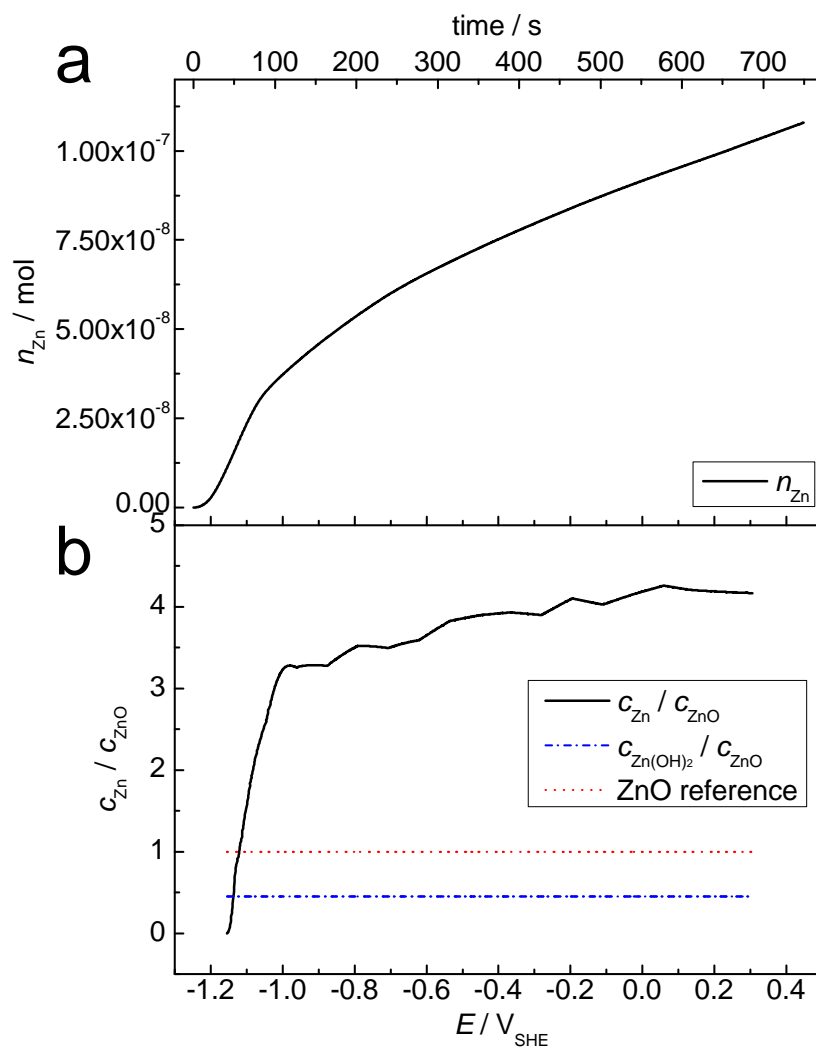


Figure 3: (a) The amount n_{Zn} of Zn which is oxidised to Zn^{2+} ; (b) $c_{\text{Zn}}/c_{\text{ZnO}}$ during the reversed positive scan between -1.15 V and 0.3 V in the CV shown in Fig. 2. The dash-dot line in (b) shows $c_{\text{Zn(OH)}_2}/c_{\text{ZnO}}$.

308 of $\text{Im}(J_1)$ is different. At wavelength of 700 nm and 650 nm, $\text{Im}(J_1)$ has the
309 same magnitude as as $\text{Re}(J_1)$. The agreement between imaginary and real
310 part in this wavelength region confirms the assumption that the imaginary
311 part in this region of the spectrum is governed by the substrate and the
312 films are non-absorbing. For lower wavelengths, absorption in the thin oxide
313 films leads to differences, e.g. at 350 nm, at a photon energy above the ZnO
314 bandgap, $\text{Im}(J_1)$ is close to zero. Therefore, these wavelengths have not been
315 used in the analysis of the thickness, but they can be used to analyse changes
316 in the electronic structure of the oxide films.

317 Changes in the electronic structure are directly manifested in the ellip-
318 sometric parameter Ψ . Fig. 5 shows the lower wavelength region of spectra
319 of Ψ during the CV. The bending at ≈ 380 nm is a typical feature for the
320 formation of zinc oxide, as it is related to the main electronic absorption of
321 ZnO, which, because of the surface selection rule, affects the amplitude re-
322 flection coefficient in p-polarisation, but not in s-polarisation, hence leading
323 to an absorption feature in the spectrum [28]. No film absorption is visi-
324 ble for the initially mechanically polished Zn substrate [7]. The first curve
325 (front black) is Ψ prior to the application of an electrochemical potential.
326 The bending at 380 nm is not observed. After switching E to +0.3 V the
327 measured Ψ changes immediately, the typical bend at 380 nm is observed
328 (second curve in Fig. 5), and remains present while scanning in cathodic
329 direction. During negative scanning, the feature becomes more pronounced
330 and disappears at strongly cathodic potentials. The apparent contradiction
331 that the absorption feature increases while the layer thickness decreases can
332 be understood as the superposition of a potential-dependent phenomenon,

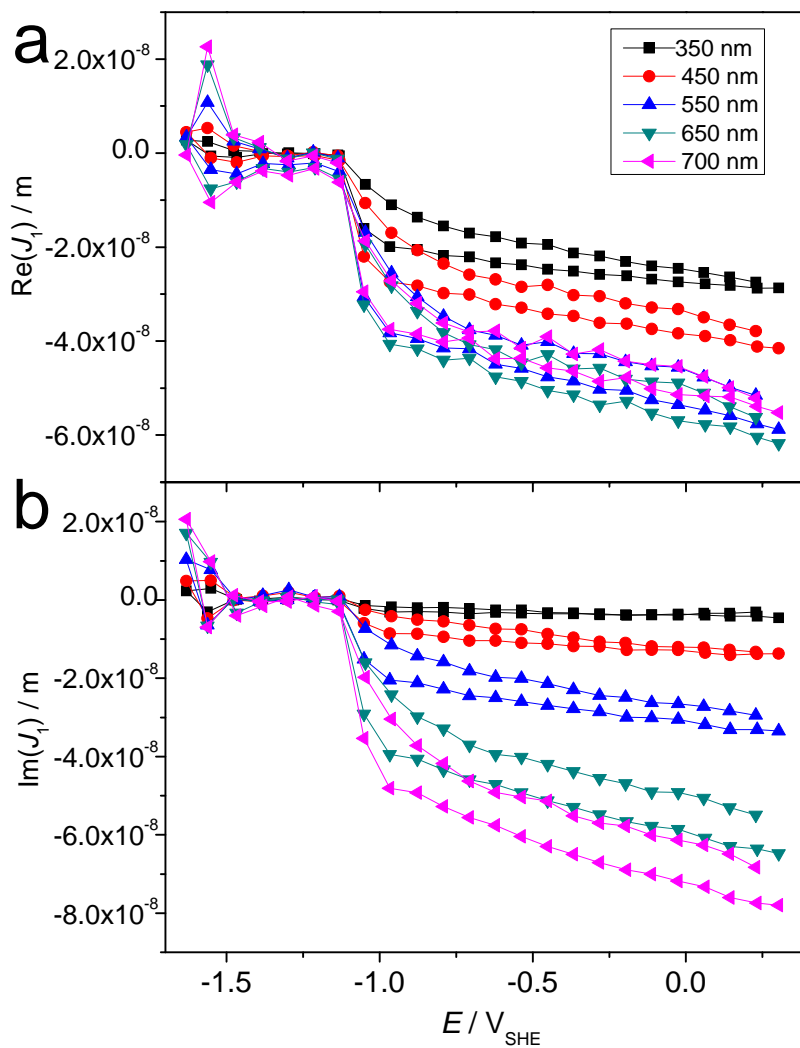


Figure 4: Real (a) and imaginary (b) part of J_1 at different wavelengths as indicated in the graph for different potentials during the CV shown in Fig. 2.

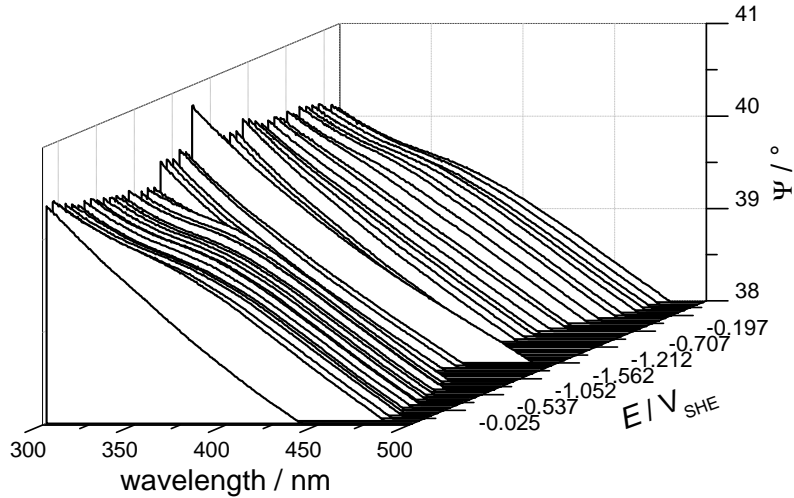


Figure 5: Spectra of Ψ for different potentials during the CV.

333 i.e. the reduction of oxide from the ZnO layer, with a time-dependent phe-
 334 nomenon, the ageing of the oxidic film. As shown in Fig. 2, d jumps from 0.8
 335 nm to 2.2 nm on the application of an anodic potential, with a subsequent
 336 gradual decrease during negative scanning. In the potential window of the
 337 cathodic reduction peak C_2 , the thickness decreased from more than 2 nm
 338 to almost zero, indicating the disappearance of oxide. In the same potential
 339 window, the absorption <380 nm in Ψ also disappeared. During the reversed
 340 positive scanning, the bend in Ψ at 380 nm was not observed until the poten-
 341 tial reached -0.6 V, subsequently becoming more pronounced with further
 342 potential increase. The absence of the absorption feature in the spectra of Ψ
 343 at potentials below -1.1 V is a good indicator that oxide is indeed absent in
 344 this potential regime, which confirms previous works [9].

345 The fact that there is only a small increase of c_{Zn} during the phase of linear

346 thickness growth, and at the same time the ZnO absorption in the spectrum
347 becomes more pronounced while in the reductive scan also an increasing
348 prominence of the oxidic absorption is observed leads to the following model.
349 Upon oxidation, $\text{Zn}(\text{OH})_2$ is initially formed on the surface. This hydrox-
350 ide ages transforming into ZnO, while the layer thickness remains constant
351 during this process. Newly formed species will always be hydroxide, which
352 ages at a certain rate, keeping the overall density constant with increasing
353 thickness.

354 *3.3. In situ SE and CA during potentiostatic oxide growth*

355 In the CV experiments, the separation of time-dependent and potential-
356 dependent features is not always possible. To have a more precise control, the
357 growth kinetics of the interfacial film on Zn during a potentiostatic process
358 was investigated. Fig. 6 shows the applied potential (a) and the correspond-
359 ing current (b) from a CA experiment. The sample was first reduced at -1.3
360 V for 60 min, then anodically oxidised initially at $+0.2$ V and subsequently
361 at $+1.2$ V for 60 min each. At the potential of -1.3 V (Fig. 1 and 2), the
362 oxide present on the surface of Zn is reduced. During the cathodic reduction
363 and further anodic oxidation processes, in situ SE was measured to determine
364 the behaviour of the oxide film. Fig. 6c indicates the thickness calculated
365 through the perturbation approach (section 2.5). Upon application of the ca-
366 thodic potential (Fig. 6b) the cathodic current diminishes fast, and after 30 s,
367 the current reaches constant value close to 0, remaining constant throughout
368 this potential stage.

369 The thickness of the film (Fig. 6c) follows closely the applied potential.
370 Because of the low speed of each SE measurement (>30 s) and the fast re-

371 duction process, the first measurement was complete after the fast reduction
372 process finished. The film thickness during the complete reduction period
373 was close to zero. The noise in the thickness data between 15 and 30 min
374 measurement time is tentatively attributed to entrapment of a gas bubble in
375 the optical path.

376 Zn is oxidised at potentials above -0.5 V. Fig. 6b shows the current
377 response after the applied potential is changed from -1.3 V to $+0.2$ V and
378 subsequently from $+0.2$ to $+1.2$ V. The current density changes to an anodic
379 maximum of $\approx 160 \mu\text{A}/\text{cm}^{-2}$ in <1 s. Subsequently, the current decreases
380 within ≈ 1 min to less than $15 \mu\text{A cm}^{-2}$. At the same time, the thickness of
381 the oxide film increases from almost zero to ≈ 1.9 nm. The increased noise in
382 the thickness results around the first potential jump originate from the fact
383 that the resulting values of J_1 are in a region of extremely high sensitivity
384 to layer parameters. Here, the data analysis procedure does hardly provide
385 stable results, which is reflected in extremely large errors of these values,
386 which are not shown as they would dominate the plot. In regions with stable
387 results (times before 60 min and after 75 min), the statistical uncertainty
388 of the thickness is ≈ 0.2 nm. After the fast initial current decrease, the
389 current decreases slowly with time, as shown in detail in the inset in Fig.
390 6b. The anodic current originates from the oxidation process of Zn on the
391 surface, and the non-zero value of the current indicates a slow but continuous
392 growth of oxide on the electrode surface. In the region where the analysis
393 procedure yields stable results, the thickness for higher times indeed shows
394 a slow increase with time. A closer look into the current response (Fig. 6b
395 inset) also indicates a current decrease after ≈ 30 min. Afterwards, $I \approx 2 - 3$

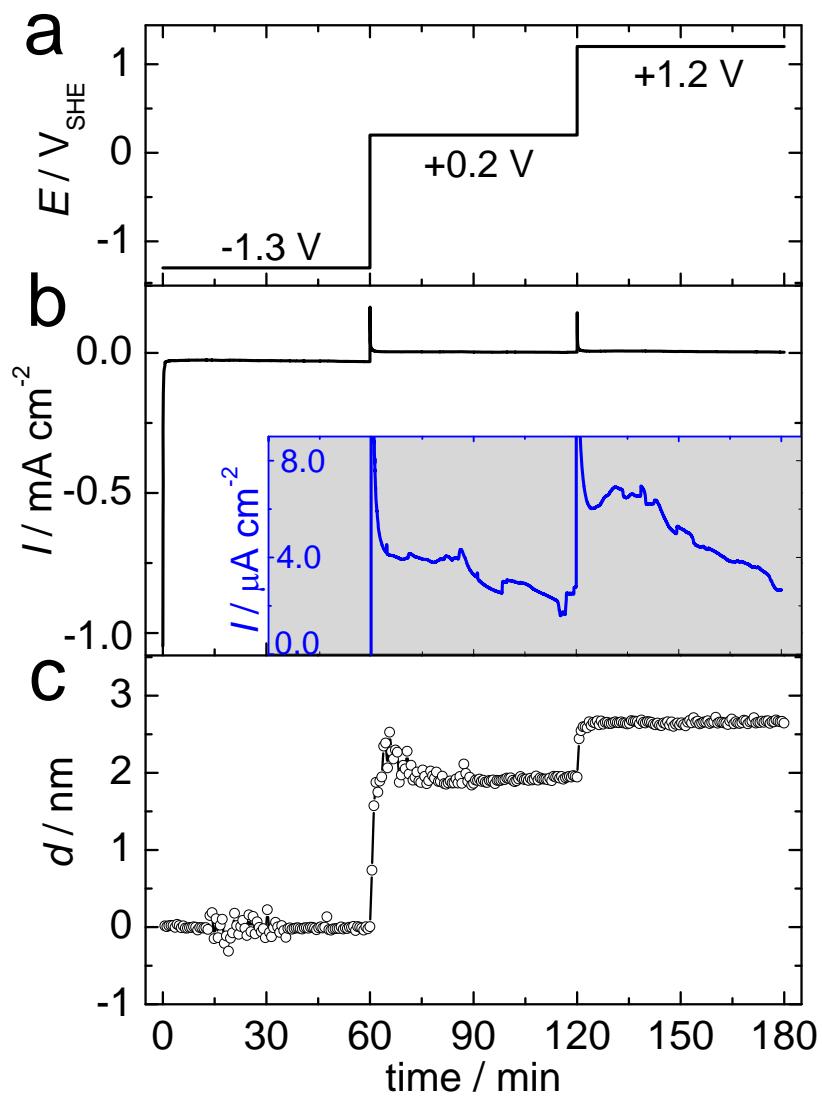


Figure 6: (a) Applied potential profile, 0–60 min: -1.3 V, 60–120 min: $+0.2$ V, 120–180 min: $+1.2$ V; (b) corresponding current densities (inset: enlarged plot of current density from 60 min to 120 min) and (c) thickness as a function of time. Errors of the thickness are ≈ 0.2 nm

396 $\mu\text{A cm}^{-2}$.

397 A further increase in the thickness of the oxide film is observed when the
398 potential is increased from +0.2 V to +1.2 V. As observed during the first
399 jump, the layer thickness increases is fast during the potential jump to 2.4
400 nm, and a further slow growth leads after 1 h to a value of 2.7 nm.

401 An integration of the current as explained in Sec. 3.2 (eq. 10) and determi-
402 nation of c_{Zn} (11) was performed for the CA data. The results are displayed
403 in Fig. 7. Comparing the increase in n_{Zn} to the increase in thickness (Fig.
404 6c) shows major differences in the time-evolution of the curves. The initial
405 current spike contains only a relatively small ($\approx 10\%$) fraction of the totally
406 transferred charge, but almost the complete growth in thickness. (In the
407 analysis here, the possibility of having missed a relatively large portion of
408 the current due to the limited bandwidth of the potentiostat is disregarded.)
409 The ratio $c_{\text{Zn}}/c_{\text{ZnO}}$ reaches 0.5 within the first 30 s after the potential jump,
410 and increases above 1 after 5 min. These results indicate that metal disso-
411 lution into the electrolyte is occurring at these potentials through the oxide
412 layer, which does not grow in size. During the second potential jump, the
413 ratio $c_{\text{Zn}}/c_{\text{ZnO}}$ decreases initially, although n_{Zn} increases. This result can be
414 easily explained by the formation of another layer on the surface of the first
415 layer. This layer experiences a subsequent densification.

416 Fig. 8 shows the plot of real and imaginary part of J_1 at different applied
417 potentials for different wavelengths. The time evolution of these curves leads
418 to the same conclusions as drawn for the CV measurements, namely that
419 absorption of the film changes during the potential treatment.

420 Fig. 9a shows the data for Ψ from the SE measurements during CA. The

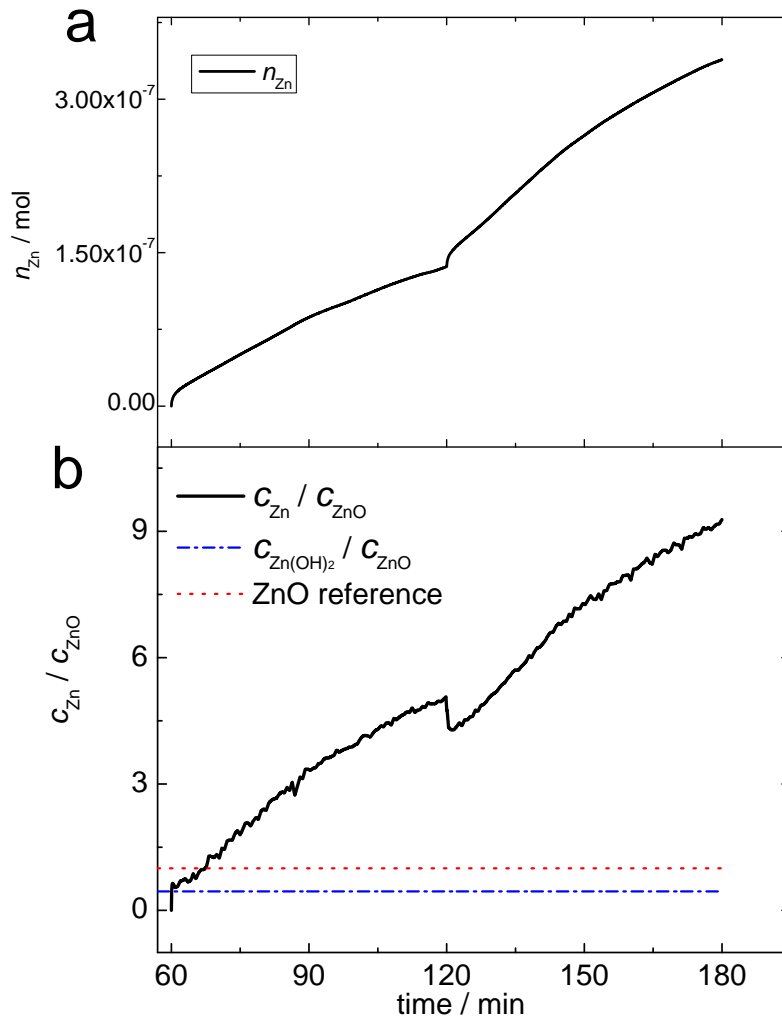


Figure 7: (a) Amount n_{Zn} of Zn which is oxidised to Zn^{2+} ; and (b) the ratio between c_{Zn} and c_{ZnO} during CA. The dash-dot line in (b) shows $c_{\text{Zn(OH)}_2} / c_{\text{ZnO}}$.

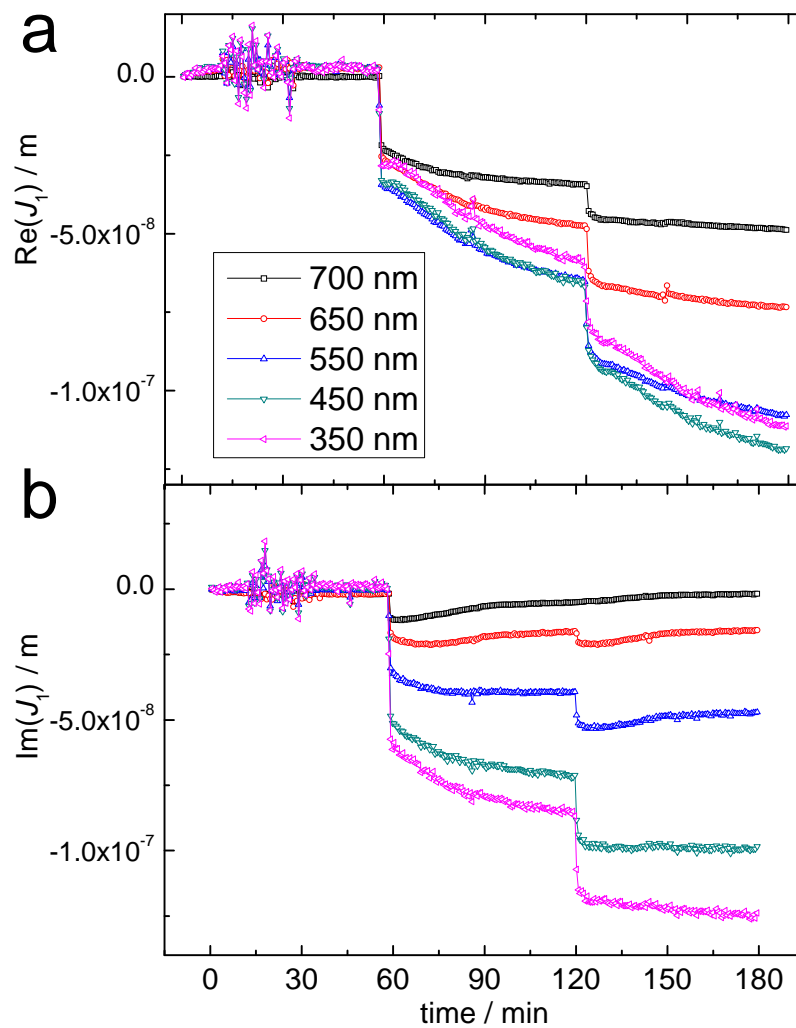


Figure 8: Real (a) and imaginary (b) part of J_1 at different wavelengths as indicated in the graph during the CA measurements shown in Fig. 6.

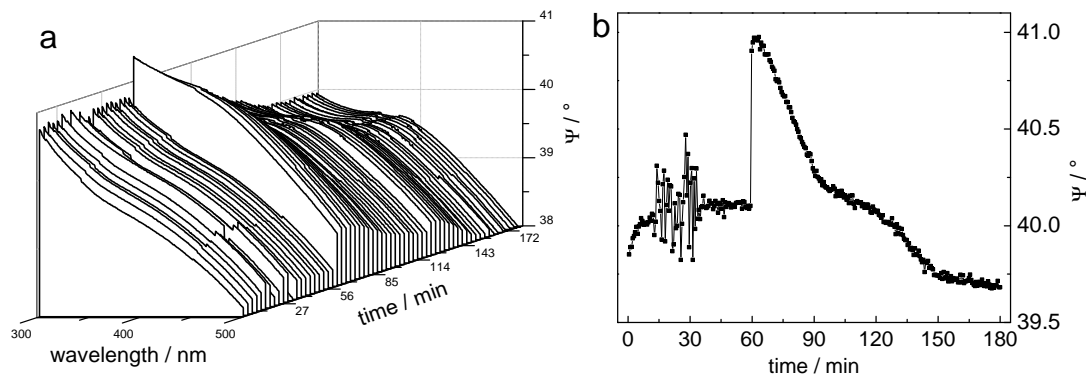


Figure 9: (a) Plot of the measured Ψ with time during the chronoamperometry. (b) Ψ at wavelength of 350 nm during the chronoamperometry

421 bending of Ψ at 380 nm is absent at a potential of -1.3 V. After switching
 422 the potential from -1.3 to $+0.2$ V, Ψ changes immediately, and the bending
 423 of Ψ at 380 nm, which indicates the presence of an absorption, becomes more
 424 and more pronounced. Fig. 9b shows the evolution with time of Ψ at 350 nm
 425 as an example of the overall curve, which also shows changes within 30 min
 426 after the respective potential jump. This behaviour indicates the continuous
 427 transformation of the layer into a “mature” oxide layer with only a small
 428 increase in thickness of the layer. After both potential jumps, this change in
 429 the absorption spectrum of the layer stops after 30 min, after which curve
 430 enters a second phase. The same second phase is also observed in the inte-
 431 grated currents in Fig. 7, which continues to increase with a slightly lower
 432 slope. The behaviour indicates that an initially formed layer is transform-
 433 ing gradually into a layer of ZnO. In the second phase, metal dissolution
 434 is still occurring, at a constant absorption spectrum (and hence, electronic
 435 structure) of the oxide layer.

436 4. Summary and Conclusions

437 During both potentiostatic treatment of zinc as well as CV measurements
438 in alkaline carbonate solution, an analysis of ellipsometric data on the basis
439 of Lekner's perturbation parameter enables a direct determination of the
440 layer thickness of the interfacial region without the need of an optical model.
441 Additionally, the analysis of the ellipsometric parameter Ψ yields insight into
442 the change in the absorption spectra of the interfacial region.

443 During dynamic potential scanning, the thinning of the oxide during the
444 negative scanning and subsequent growth of oxide in the reversed positive
445 scanning on Zn can be followed. CV-like plots in the format of thickness,
446 real or imaginary part of the perturbation parameter instead of current as
447 the function of potential are obtained. The thickness of the interfacial film,
448 which consists mainly of ZnO, grows linearly with potential at potentials
449 above the potential of the A_2 current peak. This behaviour is found for
450 passive materials as well. The characteristic electronic absorption of ZnO is
451 present in the films in this potential regime.

452 Results of potential jump experiments show that the interfacial oxide
453 layer grows in thickness rapidly upon the application of anodic potentials
454 within the time scale of less than 1 min, and the increasing of the thickness
455 during the subsequent 59 min is lower than that during the first minute.
456 The electronic absorption of ZnO occurs quickly and subsequently takes a
457 certain time (≈ 30 min) to fully evolve in the spectra, showing an ageing of
458 the interfacial layer. The analysis of the concentration of Zn^{2+} in the layer
459 indicates an ongoing dissolution of Zn into the electrolyte. The rate of Zn
460 dissolution needs to be measured independently by analysis of the electrolyte

461 to verify that the residual current is related mainly to metal dissolution.

462 More information about the electronic structure of the interfacial films are
463 expected to be obtained from a more detailed analysis of $\epsilon_l(\lambda)$. It is especially
464 promising to analyse the changes in the spectra of a dynamic system, i.e.
465 while Zn transport during the layer is ongoing, as is the case here in during the
466 CA measurements. The ageing observed here may be related to the ongoing
467 transport of zinc through the layer. A comparison to passive materials with
468 very low residual currents of dissolution should allow conclusions about this
469 behaviour. With increased time resolution of the ellipsometric experiments,
470 more details about the initial stages of the layer formation will be obtained.

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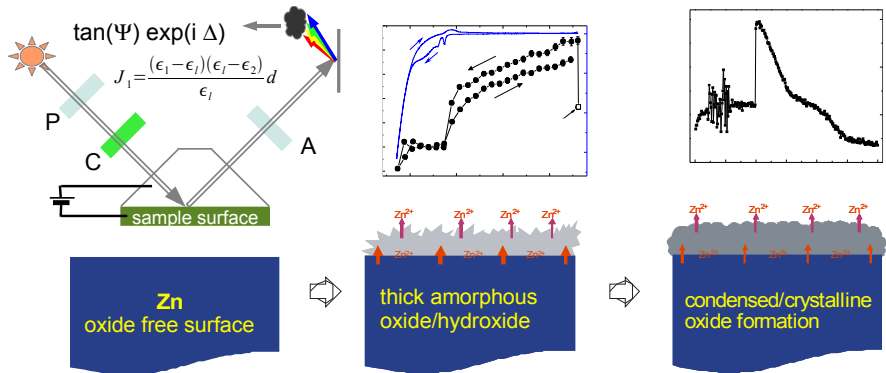
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541 **Highlights**

- 542 • thickness determination from ellipsometry without optical model
- 543 • ZnO thickness increases linearly with potential in the “pseudo-passive”
544 regime
- 545 • density of interfacial layer is constant during potential increase
- 546 • potentiostatic treatment shows an aging of the interfacial layer to yield
547 ZnO
- 548 • absorption spectrum of the oxide is modified by metal transport

549 Graphical abstract



550